Bihamiltonian Structure of the Two-component Kadomtsev-Petviashvili Hierarchy of type B

Chao-Zhong Wu
* Dingdian Xu^{\dagger} Department of Mathematical Sciences, Tsinghua University, Beijing 100084, P. R. China

Abstract

We employ a Lax pair representation of the two-component BKP hierarchy and construct its bihamiltonian structure with R-matrix techniques.

Key words: BKP hierarchy, Hamiltonian structure, R-matrix

1 Introduction

The Kadomtsev-Petviashvili (KP) hierarchy of type B (BKP for short) was introduced in [6, 7], and generalized to multi-component cases by Date, Jimbo, Kashiwara, Miwa [4] in the form of bilinear equations. Among these multi-component integrable systems, the two-component BKP hierarchy is of special interest.

In the original definition of the two-component BKP hierarchy, the solution space of tau functions can be regarded as the vacuum orbit in the two-component neutral free fermionic Fock representation of the infinite dimensional Lie algebra D_{∞} [5, 14], which corresponds to the infinite Dynkin diagram of type D [15]. The Lie algebra D_{∞} can be reduced to the affine Lie algebra $D_n^{(1)}$ under the so-called (2n-2,2)-reduction in [5], see also [14, 17]. This reduction reduces the two-component BKP hierarchy to a hierarchy that is equivalent with the Kac-Wakimoto hierarchy corresponding to the principal vertex operator realization of the basic representation of $D_n^{(1)}$, the Drinfeld-Sokolov hierarchy associated to the Lie algebra $D_n^{(1)}$ and the zeroth vertex c_0 of its Dynkin diagram, as well as the Givental-Milanov hierarchy satisfied by the total descendant for the D_n singularity,

^{*}wucz05@mails.tsinghua.edu.cn

 $^{^{\}dagger}$ xudd06@mails.tsinghua.edu.cn

see [9, 12, 13, 16, 19, 26] and references therein. Such a reduction is analogous to the one that reduces the KP hierarchy to the nth Gelfand-Dickey hierarchy (see e.g. [8]) that corresponds to the reduction of Lie algebras: $A_{\infty} \mapsto A_n^{(1)}$. So in this sense to compare the two-component BKP hierarchy with the KP hierarchy would deepen our understanding of integrable hierarchies and relevant theories, such as Jacobi/Prym varieties in algebraic geometry and Landau-Ginzburg Models of topological strings, see e.g. [22, 23, 24].

In this article our aim is to study the two-component BKP hierarchy from the view point of Hamiltonian structures. To our best knowledge, this topic has not been considered in the literature, possibly for the reason that the KP-analogue Lax pair representation of the two-component BKP hierarchy was unknown. Recall that the two-component BKP hierarchy was defined to be the bilinear equation of a single tau function:

$$\operatorname{res}_{z} z^{-1} X(\mathbf{t}; z) \tau(\mathbf{t}, \hat{\mathbf{t}}) X(\mathbf{t}'; -z) \tau(\mathbf{t}', \hat{\mathbf{t}}')$$

$$= \operatorname{res}_{z} z^{-1} X(\hat{\mathbf{t}}; z) \tau(\mathbf{t}, \hat{\mathbf{t}}) X(\hat{\mathbf{t}}'; -z) \tau(\mathbf{t}', \hat{\mathbf{t}}'), \tag{1.1}$$

where $\mathbf{t} = (t_1, t_3, t_5, \cdots)$, $\hat{\mathbf{t}} = (\hat{t}_1, \hat{t}_3, \hat{t}_5, \cdots)$, and X is a vertex operator given by

$$X(\mathbf{t};z) = \exp\left(\sum_{k \in \mathbb{Z}_+^{\text{odd}}} t_k z^k\right) \exp\left(-\sum_{k \in \mathbb{Z}_+^{\text{odd}}} \frac{2}{k z^k} \frac{\partial}{\partial t_k}\right).$$

Here the residue of a Laurent series is taken as $\operatorname{res}_z(\sum_{i\in\mathbb{Z}} f_i z^i) = f_{-1}$. In [22] Shiota proposed a scalar Lax representation of the hierarchy (1.1), though this did not attract much attention as it contains pseudo-differential operators with derivations of two spatial variables. Recently, a Lax pair representation of the two-component BKP hierarchy was found by Liu, Zhang and one of the authors [19]. It was shown that the hierarchy (1.1) can be redefined by certain extension of the following Lax equations (see Section 3 below):

$$\frac{\partial P}{\partial t_k} = [(P^k)_+, P], \quad \frac{\partial \hat{P}}{\partial t_k} = [(P^k)_+, \hat{P}], \tag{1.2}$$

$$\frac{\partial P}{\partial \hat{t}_k} = [-(\hat{P}^k)_-, P], \quad \frac{\partial \hat{P}}{\partial \hat{t}_k} = [-(\hat{P}^k)_-, \hat{P}] \tag{1.3}$$

with $k \in \mathbb{Z}_+^{\text{odd}}$, where

$$P = D + \sum_{i>1} u_i D^{-i}, \quad \hat{P} = D^{-1} \hat{u}_{-1} + \sum_{i>1} \hat{u}_i D^i \text{ with } D = \frac{\mathrm{d}}{\mathrm{d}x}$$

are pseudo-differential operators such that $P^* = -DPD^{-1}$, $\hat{P}^* = -D\hat{P}D^{-1}$. Note that the first equation in (1.2) is just the Lax formulation of the BKP hierarchy appearing in [6]. Our arguments will be based on the Lax pair representation (1.2), (1.3) of the two-component BKP hierarchy.

Observe that the expression (1.2), (1.3) is similar to the Lax pair representation of the two-dimensional Toda hierarchy [25], which carries a tri-Hamiltonian structure [1]. Following the idea of [1], we want to use the R-matrix theory to construct Hamiltonian structures of the two-component BKP hierarchy (1.2), (1.3).

We are also motivated by the recent work [2], in which Carlet, Dubrovin and Mertens constructed an infinite-dimensional Frobenius manifold underlying the two-dimensional Toda hierarchy. Due to the similarity of the Lax representations mentioned above, we expect that there also exists an infinite dimensional Frobenius manifold that underlies the two-component BKP hierarchy. A hint is that the potential F (in the notion of [23], namely the dispersionless limit of the logarithm of the tau function, see Section 3 below) of the dispersionless two-component BKP hierarchy was discovered to satisfy certain infinite-dimensional WDVV-type associativity equation [3]. While in the finite-dimensional case, the concept of Frobenius manifolds [10] is known as a geometric description of the WDVV equations, and associated to certain nondegenerate Frobenius manifold there lies a Poisson pencil so that a bihamiltonian hierarchy can be constructed [11]. We hope that this article and follow-up work might help to understand the theory of infinite-dimensional manifolds.

This article is arranged as follows. In next section we recall the definition and some properties of pseudo-differential operators introduced in [19], and in Section 3 we recall the Lax pair representation of the two-component BKP hierarchy. In Sections 4 and 5, an *R*-matrix will be used to construct Poisson brackets on an algebra of pseudo-differential operators, and then after appropriate reductions of the Poisson brackets we obtain a bihamiltonian structure of the two-component BKP hierarchy. In Section 6 we compute the dispersionless limit of this bihamiltonian structure. Finally some remarks are given in Section 7.

2 Pseudo-differential operators

For preparation we recall the notion of pseudo-differential operators over a ring with certain gradation as introduced in [19].

Let \mathcal{A} be a ring, and $D: \mathcal{A} \to \mathcal{A}$ be a derivation. The algebra of usual

pseudo-differential operators is

$$\mathcal{D}^{-} = \left\{ \sum_{i < \infty} f_i D^i \mid f_i \in \mathcal{A} \right\}. \tag{2.1}$$

This algebra is topologically complete with a topological basis given by the following filtration:

$$\cdots \subset \mathcal{D}_{(d-1)}^{-} \subset \mathcal{D}_{(d)}^{-} \subset \mathcal{D}_{(d+1)}^{-} \subset \cdots, \quad \mathcal{D}_{(d)}^{-} = \left\{ \sum_{i \leq d} f_i D^i \mid f_i \in \mathcal{A} \right\},$$

and in this algebra two elements are multiplied as series of the following product of monomials:

$$fD^{i} \cdot gD^{j} = \sum_{r>0} {i \choose r} f D^{r}(g) D^{i+j-r}, \quad f, g \in \mathcal{A}.$$
 (2.2)

Assume there is a gradation on A such that

$$\mathcal{A} = \prod_{i>0} \mathcal{A}_i, \quad D: \mathcal{A}_i \to \mathcal{A}_{i+1}, \quad \mathcal{A}_i \cdot \mathcal{A}_j \subset \mathcal{A}_{i+j},$$

and consider the linear space

$$\mathcal{D} = \left\{ \sum_{i \in \mathbb{Z}} f_i D^i \mid f_i \in \mathcal{A}
ight\}.$$

Obviously $\mathcal{D}^- \subset \mathcal{D}$.

For any $k \in \mathbb{Z}$, denote by \mathcal{D}_k the set of homogeneous operators with degree k in \mathcal{D}^- , i.e.,

$$\mathcal{D}_k = \left\{ \sum_{i \le k} f_i D^i \mid f_i \in \mathcal{A}_{k-i} \right\}.$$

Let \mathcal{D}^+ be a subspace of \mathcal{D} that reads

$$\mathcal{D}^{+} = \bigcup_{d \in \mathbb{Z}} \mathcal{D}_{(d)}^{+}, \quad \mathcal{D}_{(d)}^{+} = \prod_{k \ge d} \mathcal{D}_{k}, \tag{2.3}$$

and \mathcal{D}^+ have a topological basis given by the filtration

$$\cdots \supset \mathcal{D}^+_{(d-1)} \supset \mathcal{D}^+_{(d)} \supset \mathcal{D}^+_{(d+1)} \supset \cdots$$

In fact, every element $A \in \mathcal{D}^+$ has the following normal expansion [19]

$$A = \sum_{i \in \mathbb{Z}} \left(\sum_{j \ge \max\{0, m-i\}} a_{i,j} \right) D^i, \quad a_{i,j} \in \mathcal{A}_j$$

with some integer m. Note that $\mathcal{D}_k \cdot \mathcal{D}_l \subset \mathcal{D}_{k+l}$ according to the multiplication defined by (2.2), then this multiplication can be naturally extended to \mathcal{D}^+ such that \mathcal{D}^+ becomes an associative algebra.

Definition 2.1 ([19]) Elements of \mathcal{D}^- (resp. \mathcal{D}^+) are called pseudo-differential operators of the first type (resp. the second type) over \mathcal{A} . The intersection of \mathcal{D}^- and \mathcal{D}^+ in \mathcal{D} is denoted by

$$\mathcal{D}^b = \mathcal{D}^- \cap \mathcal{D}^+.$$

and its elements are called bounded pseudo-differential operators.

Sometimes to indicate the ring A and the derivation D, we will use the notations $\mathcal{D}^{\pm}(A, D)$ instead of \mathcal{D}^{\pm} .

Pseudo-differential operators of the second type have similar properties to those of the operators in \mathcal{D}^- . For any operator

$$A = \sum_{i \in \mathbb{Z}} f_i D^i \in \mathcal{D}^{\pm}, \tag{2.4}$$

its positive part, negative part, residue and adjoint operator are defined to be respectively

$$A_{+} = \sum_{i>0} f_i D^i, \quad A_{-} = \sum_{i<0} f_i D^i,$$
 (2.5)

res
$$A = f_{-1}, \quad A^* = \sum_{i \in \mathbb{Z}} (-D)^i \cdot f_i.$$
 (2.6)

Note that the formulae (2.5) give two projections of \mathcal{D} , and they induce the following decompositions of spaces

$$\mathcal{D}^{\pm} = (\mathcal{D}^{\pm})_{+} \oplus (\mathcal{D}^{\pm})_{-}. \tag{2.7}$$

Particularly one sees that

$$(\mathcal{D}^-)_+ \subset \mathcal{D}^b, \quad (\mathcal{D}^+)_- \subset \mathcal{D}^b.$$
 (2.8)

An element A of $(\mathcal{D}^{\pm})_+$ is called a differential operator. Let A(f) denote the action of a differential operator A on $f \in \mathcal{A}$.

Elements of the quotient space $\mathcal{F} = \mathcal{A}/(D(\mathcal{A}) \oplus \mathbb{C})$ are called *local functionals*, which are denoted as

$$\int f \, \mathrm{d}x = f + D(\mathcal{A}), \quad f \in \mathcal{A}.$$

Introduce a map

$$\langle \rangle : \mathcal{D} \to \mathcal{F}, \quad A \mapsto \langle A \rangle = \int \operatorname{res} A \, \mathrm{d}x.$$
 (2.9)

Then the pairing

$$\langle A, B \rangle = \langle AB \rangle \tag{2.10}$$

defines an inner product on each of \mathcal{D}^{\pm} .

Given any subspace $\mathcal{S} \subset \mathcal{D}^{\pm}$, we denote by \mathcal{S}^* the dual space of \mathcal{S} (c.f. the notation of adjoint operators). Via the above inner product, we have the following identification of dual spaces

$$(\mathcal{D}^{\pm})^* = \mathcal{D}^{\pm}. \tag{2.11}$$

Consider the decompositions (2.7), it is easy to see that

$$\left((\mathcal{D}^{\pm})_{\pm}\right)^* = (\mathcal{D}^{\pm})_{\mp}.$$

We also decompose \mathcal{D}^{\pm} as

$$\mathcal{D}^{\pm} = \mathcal{D}_0^{\pm} \oplus \mathcal{D}_1^{\pm}, \tag{2.12}$$

where

$$\mathcal{D}_{\nu}^{\pm} = \left\{ A \in \mathcal{D}^{\pm} \mid A^* = (-1)^{\nu} A \right\}, \quad \nu = 0, 1.$$

Since $\langle A \rangle = -\langle A^* \rangle$ for any $A \in \mathcal{D}^{\pm}$, then the dual subspaces of \mathcal{D}^{\pm}_{ν} read

$$(\mathcal{D}_{\nu}^{\pm})^* = \mathcal{D}_{1-\nu}^{\pm}, \quad \nu = 0, 1.$$
 (2.13)

For more details on properties of pseudo-differential operators one can refer to [8, 19].

3 The two-component BKP hierarchy

The two types of pseudo-differential operators serve in [19] to give a scalar Lax pair representation of the two-component BKP hierarchy, which is reviewed as follows.

Let \tilde{M} be an infinite-dimensional manifold with local coordinates

$$(a_1, a_3, a_5, \ldots, b_1, b_3, b_5, \ldots),$$

and $\tilde{\mathcal{A}}$ be the algebra of differential polynomials on \tilde{M} :

$$\tilde{\mathcal{A}} = C^{\infty}(\tilde{M})[[a_k^{(s)}, b_k^{(s)} \mid k \in \mathbb{Z}_+^{\mathrm{odd}}, s \ge 1]].$$

We assign a gradation on $\tilde{\mathcal{A}}$ by

$$\deg f = 0 \text{ for } f \in C^{\infty}(\tilde{M}), \quad \deg a_k^{(s)} = \deg b_k^{(s)} = s$$

which make $\tilde{\mathcal{A}}$ a topologically complete algebra:

$$\tilde{\mathcal{A}} = \prod_{i>0} \tilde{\mathcal{A}}_i, \quad \tilde{\mathcal{A}}_i \cdot \tilde{\mathcal{A}}_j \subset \tilde{\mathcal{A}}_{i+j}.$$

Note that this gradation is induced from the derivation

$$D: \tilde{\mathcal{A}} \to \tilde{\mathcal{A}}, \quad D = \sum_{s \geq 0} \sum_{k \in \mathbb{Z}_{\perp}^{\mathrm{odd}}} \left(a_k^{(s+1)} \frac{\partial}{\partial a_k^{(s)}} + b_k^{(s+1)} \frac{\partial}{\partial b_k^{(s)}} \right)$$

with $a_k^{(0)} = a_k$, $b_k^{(0)} = b_k$. So one can define the algebras $\tilde{\mathcal{D}}^{\pm} = \mathcal{D}^{\pm}(\tilde{\mathcal{A}}, D)$ of pseudo-differential operators as was done in last section.

Introduce two operators

$$\Phi = 1 + \sum_{i>1} a_i D^{-i} \in \tilde{\mathcal{D}}^-, \quad \Psi = 1 + \sum_{i>1} b_i D^i \in \tilde{\mathcal{D}}^+, \tag{3.1}$$

where $a_2, a_4, a_6, \ldots, b_2, b_4, b_6, \cdots \in \tilde{\mathcal{A}}$ are determined by the following conditions

$$\Phi^* = D\Phi^{-1}D^{-1}, \quad \Psi^* = D\Psi^{-1}D^{-1}. \tag{3.2}$$

Then the two-component BKP hierarchy (1.1) can be redefined to be

$$\frac{\partial \Phi}{\partial t_k} = -(P^k)_- \Phi, \quad \frac{\partial \Psi}{\partial t_k} = \left((P^k)_+ - \delta_{k1} \hat{P}^{-1} \right) \Psi, \tag{3.3}$$

$$\frac{\partial \Phi}{\partial \hat{t}_k} = -(\hat{P}^k)_- \Phi, \quad \frac{\partial \Psi}{\partial \hat{t}_k} = (\hat{P}^k)_+ \Psi, \tag{3.4}$$

where $k \in \mathbb{Z}_{+}^{\text{odd}}$, and the operators P, \hat{P} read

$$P = \Phi D \Phi^{-1} \in \tilde{\mathcal{D}}^-, \quad \hat{P} = \Psi D^{-1} \Psi^{-1} \in \tilde{\mathcal{D}}^+.$$
 (3.5)

The operators P, \hat{P} have the following expressions:

$$P = D + \sum_{i \ge 1} u_i D^{-i}, \quad \hat{P} = D^{-1} \hat{u}_{-1} + \sum_{i \ge 1} \hat{u}_i D^i, \tag{3.6}$$

with $\hat{u}_{-1} = (\Psi^{-1})^*(1)$, and they satisfy

$$P^* = -DPD^{-1}, \quad \hat{P}^* = -D\hat{P}D^{-1}, \tag{3.7}$$

which implies

$$(P^k)_+(1) = 0, \quad (\hat{P}^k)_+(1) = 0, \quad k \in \mathbb{Z}_+^{\text{odd}}.$$
 (3.8)

Observe that the coefficients of P and \hat{P} are elements of the algebra $\tilde{\mathcal{A}}$, and among these coefficients the ones with odd subscript are independent, while the others are determined by the conditions (3.7). Assume that

$$\mathbf{u} = (u_1, u_3, \dots, \hat{u}_{-1}, \hat{u}_1, \hat{u}_3, \dots)$$
(3.9)

serves as a coordinate of some infinite-dimensional manifold M, then the algebra \mathcal{A} of differential polynomials on M reads

$$\mathcal{A} = C^{\infty}(M)[[\mathbf{u}^{(s)} \mid s \ge 1]],$$

which is a subalgebra of $\tilde{\mathcal{A}}$. Similarly as above, one can assign a gradation to \mathcal{A} that is induced from the derivation

$$D: \mathcal{A} \to \mathcal{A}, \quad D = \sum_{s>0} \mathbf{u}^{(s+1)} \cdot \frac{\partial}{\partial \mathbf{u}^{(s)}}$$

with $\mathbf{u}^{(0)} = \mathbf{u}$, and then define the algebras $\mathcal{D}^{\pm} = \mathcal{D}^{\pm}(\mathcal{A}, D)$ of pseudo-differential operators over \mathcal{A} .

Clearly $P \in \mathcal{D}^-$, $\hat{P} \in \mathcal{D}^+$. When the two-component BKP hierarchy (3.3), (3.4) is restricted from $\tilde{\mathcal{A}}$ to \mathcal{A} , it becomes

$$\frac{\partial P}{\partial t_k} = [(P^k)_+, P], \quad \frac{\partial \hat{P}}{\partial t_k} = [(P^k)_+, \hat{P}], \tag{3.10}$$

$$\frac{\partial P}{\partial \hat{t}_k} = [-(\hat{P}^k)_-, P], \quad \frac{\partial \hat{P}}{\partial \hat{t}_k} = [-(\hat{P}^k)_-, \hat{P}]$$
(3.11)

with $k \in \mathbb{Z}_+^{\text{odd}}$. In the present article we regard the two-component BKP hierarchy as the evolutionary equations (3.10), (3.11) defined on the algebra A.

In fact, the hierarchy (3.10), (3.11) possesses a tau function $\tau = \tau(\mathbf{t}, \hat{\mathbf{t}})$ defined by

$$\omega = d(2 \partial_x \log \tau) \quad \text{with} \quad x = t_1,$$
 (3.12)

where ω is the following closed 1-form:

$$\omega = \sum_{k \in \mathbb{Z}_+^{\text{odd}}} (\operatorname{res} P^k dt_k + \operatorname{res} \hat{P}^k d\hat{t}_k).$$

This tau function solves the bilinear equation (1.1), which is the original definition of the two-component BKP hierarchy.

Remark 3.1 The dispersionless limit of the flows (3.10), (3.11) first exists in [23], where Takasiki also considered the dispersionless limit of the logarithm of the tau function as given in (3.12). Inspired by [23], Chen and Tu [3] discovered that the leading term of $\log \tau$ solves an infinite-dimensional associativity equation of WDVV type.

4 R-matrix and pseudo-differential operators

To show that the two-component BKP hierarchy (3.10), (3.11) possesses a bihamiltonian structure, we need to construct a Poisson pencil for it. The method is to use the standard R-matrix theory and introduce Poisson brackets on a Lie algebra (see [21, 18, 20] and references therein), then restrict the Poisson brackets to certain submanifold of the Lie algebra. Our approach is similar with that used by Carlet [1] for the two-dimensional Toda hierarchy.

We first recall the R-matrix formalism. Let \mathfrak{g} be a Lie algebra, and $R:\mathfrak{g}\to\mathfrak{g}$ be a linear transformation. Then R is called an R-matrix [21] on \mathfrak{g} if it defines a Lie bracket by

$$[X,Y]_R = [R(X),Y] + [X,R(Y)], \quad X,Y \in \mathfrak{g}.$$
 (4.1)

A sufficient condition for a transformation R being an R-matrix is that R solves the modified Yang-Baxter equation [21]

$$[R(X), R(Y)] - R([X, Y]_R) = -[X, Y]$$
(4.2)

for all $X, Y \in \mathfrak{g}$.

Assume that \mathfrak{g} is an associative algebra, with the Lie bracket defined naturally by commutators, and there is a map $\langle \ \rangle : \mathfrak{g} \to \mathbb{C}$ that defines a non-degenerate symmetric invariant bilinear form (inner product) $\langle \ , \ \rangle$ by

$$\langle X, Y \rangle = \langle XY \rangle = \langle YX \rangle, \quad X, Y \in \mathfrak{g}.$$

Via this inner product one can identify \mathfrak{g} with its dual space \mathfrak{g}^* . The tangent and the cotangent bundles of \mathfrak{g} are denoted by $T\mathfrak{g}$ and $T^*\mathfrak{g}$ respectively, with fibers $T_A\mathfrak{g} = \mathfrak{g}$ and $T_A^*\mathfrak{g} = \mathfrak{g}^*$ at every point $A \in \mathfrak{g}$.

Let R^* be the adjoint transformation of R with respect to the above inner product. We introduce the notations of the symmetric and the antisymmetric parts of R respectively as

$$R_s = \frac{1}{2}(R + R^*), \quad R_a = \frac{1}{2}(R - R^*).$$

The R-matrix formalism is briefly stated as follows. Given an R-matrix $R: \mathfrak{g} \to \mathfrak{g}$ that satisfies certain conditions, there define three compatible Poisson brackets on \mathfrak{g} , say, the linear, the quadratic and the cubic brackets in the notion of [18, 20].

In particular, let us recall the quadratic bracket, which will be used to construct a Poisson pencil for the two-component BKP hierarchy.

Lemma 4.1 ([18, 20]) Let f, g be two arbitrary smooth functions on \mathfrak{g} , and ∇f , $\nabla g \in T_A^*\mathfrak{g}$ be their gradients at any point $A \in \mathfrak{g}$. Given a linear

transformation $R: \mathfrak{g} \to \mathfrak{g}$, if both R and its anti-symmetric part R_a satisfy the modified Yang-Baxter equation (4.2), then the quadratic bracket

$$\{f,g\}(A) = \frac{1}{4} \left(\langle [A, \nabla f], R(A\nabla g + \nabla g \cdot A) \rangle - \langle [A, \nabla g], R(A\nabla f + \nabla f \cdot A) \rangle \right)$$

$$(4.3)$$

defines a Poisson bracket on g.

Note that the bracket (4.3) can be rewritten as

$$\{f, g\}(A) = \langle \nabla f, \mathcal{P}_A(\nabla g) \rangle,$$

where $\mathcal{P}: T^*\mathfrak{g} \to T\mathfrak{g}$ is a Poisson tensor given by

$$\mathcal{P}_{A}(\nabla g) = -\frac{1}{4}[A, R(A\nabla g + \nabla g \cdot A)] - \frac{1}{4}AR^{*}([A, \nabla g]) - \frac{1}{4}R^{*}([A, \nabla g])A,$$

namely

$$\mathcal{P}_A(\nabla g) = -\frac{1}{2}A(R_s(A\nabla g) + R_a(\nabla g \cdot A)) + \frac{1}{2}(R_a(A\nabla g) + R_s(\nabla g \cdot A))A.$$
(4.4)

Henceforth we take \mathfrak{g} to be the algebra

$$\mathfrak{D} = \mathcal{D}^- \times \mathcal{D}^+$$
.

where \mathcal{D}^- and \mathcal{D}^+ are the sets of pseudo-differential operators of the first type and the second type over some differential algebra \mathcal{A} as defined in Section 2. In \mathfrak{D} the elements read $\mathbf{X} = (X, \hat{X})$, and the operations are defined diagonally as

$$(X, \hat{X}) + (Y, \hat{Y}) = (X + Y, \hat{X} + \hat{Y}), \quad (X, \hat{X})(Y, \hat{Y}) = (XY, \hat{X}\hat{Y}).$$

So $\mathfrak D$ is indeed an associative algebra. Moreover, the algebra $\mathfrak D$ is equipped with an inner product define by

$$\langle (X, \hat{X}), (Y, \hat{Y}) \rangle = \langle (X, \hat{X})(Y, \hat{Y}) \rangle = \langle X, Y \rangle + \langle \hat{X}, \hat{Y} \rangle,$$

see (2.9), (2.10). Via this inner product we have the identification of dual spaces as above:

$$\mathfrak{D}^* = (\mathcal{D}^-)^* \times (\mathcal{D}^+)^* = \mathcal{D}^- \times \mathcal{D}^+ = \mathfrak{D}.$$

Inspired by [1], we introduce a linear transformation of \mathfrak{D} as follows

$$R: \mathfrak{D} \to \mathfrak{D}, \quad (X, \hat{X}) \mapsto (X_{+} - X_{-} + 2\hat{X}_{-}, \hat{X}_{-} - \hat{X}_{+} + 2X_{+}).$$
 (4.5)

Since $R = \Pi - \tilde{\Pi}$, where

$$\Pi(X,\hat{X}) = (X_+ + \hat{X}_-, \hat{X}_- + X_+), \quad \tilde{\Pi}(X,\hat{X}) = (X_- - \hat{X}_-, \hat{X}_+ - X_+)$$

are two projections of \mathfrak{D} onto its subalgebras, more exactly,

$$\Pi \mathfrak{D} = \{ (X, X) \mid X \in \mathcal{D}^b \}, \quad \tilde{\Pi} \mathfrak{D} = (\mathcal{D}^-)_- \times (\mathcal{D}^+)_+,$$

$$\Pi^2 = \Pi, \quad \tilde{\Pi}^2 = \tilde{\Pi}, \quad \tilde{\Pi} \Pi = 0 = \Pi \tilde{\Pi}, \quad \Pi + \tilde{\Pi} = \mathrm{id},$$

then transformation R satisfies the modified Yang-Baxter equation (4.2). Hence R is an R-matrix on \mathfrak{D} .

On the other hand, with respect to the inner product on $\mathfrak D$ the adjoint transformation of R reads

$$R^*: \mathfrak{D} \to \mathfrak{D}, \quad (X, \hat{X}) \mapsto (X_- - X_+ + 2\hat{X}_-, \hat{X}_+ - \hat{X}_- + 2X_+).$$

Then the symmetric and anti-symmetric parts of R are given by

$$R_s(X,\hat{X}) = 2(\hat{X}_-, X_+), \quad R_a(X,\hat{X}) = (X_+ - X_-, \hat{X}_- - \hat{X}_+).$$
 (4.6)

Observe that R_a can be expressed as the difference of two projections onto subalgebras of \mathfrak{D} , hence R_a also solves the Yang-Baxter equation (4.2). Thus the R-matrix given in (4.5) fulfills the condition of Lemma 4.1.

We regard \mathfrak{D} as an infinite-dimensional manifold, whose coordinate is given by the coefficients of the general expression of its elements

$$\mathbf{A} = \left(\sum_{i \in \mathbb{Z}} w_i D^i, \sum_{i \in \mathbb{Z}} \hat{w}_i D^i\right) \in \mathfrak{D}.$$
 (4.7)

The set \mathcal{F} of local functionals over the differential algebra \mathcal{A} (see Section 2) plays the role of $C^{\infty}(\mathfrak{g})$. For any $F = \int f \, \mathrm{d}x \in \mathcal{F}$, the variational gradient of F at \mathbf{A} given in (4.7) is defined to be

$$\frac{\delta F}{\delta \mathbf{A}} = \left(\sum_{i \in \mathbb{Z}} D^{-i-1} \frac{\delta F}{\delta w_i(x)}, \sum_{i \in \mathbb{Z}} D^{-i-1} \frac{\delta F}{\delta \hat{w}_i(x)} \right),$$

where $\delta F/\delta w(x) = \sum_{j\geq 0} (-D)^j \left(\partial f/\partial w^{(j)}\right)$. Note that $\delta F/\delta \mathbf{A}$ is not contained in $\mathfrak{D}^* = \mathfrak{D}$ in general, so to go forward we need to do some restriction.

It shall be indicated that, in this paper we only consider functionals with variational gradients lying in \mathfrak{D} . Let \mathcal{F}_0 denote the set of such functionals.

Now we can use Lemma 4.1 and the formulae (4.4), (4.6) to obtain the following result.

Lemma 4.2 Let F and G be two arbitrary functionals in \mathcal{F}_0 . On the algebra \mathfrak{D} there is a quadratic Poisson bracket

$$\{F, G\}(\mathbf{A}) = \left\langle \frac{\delta F}{\delta \mathbf{A}}, \mathcal{P}_{\mathbf{A}} \left(\frac{\delta G}{\delta \mathbf{A}} \right) \right\rangle, \quad \mathbf{A} = (A, \hat{A}) \in \mathfrak{D},$$
 (4.8)

where the Poisson tensor $\mathcal{P}: T\mathfrak{D}^* \to T\mathfrak{D}$ is defined by

$$\mathcal{P}_{(A,\hat{A})}(X,\hat{X}) = (A(XA)_{-} - (AX)_{-}A - A(\hat{A}\hat{X})_{-} + (\hat{X}\hat{A})_{-}A,$$
$$\hat{A}(\hat{X}\hat{A})_{+} - (\hat{A}\hat{X})_{+}\hat{A} - \hat{A}(AX)_{+} + (XA)_{+}\hat{A}). \tag{4.9}$$

Aiming at Hamiltonian structures of the two-component BKP hierarchy, we need to reduce the Poisson structure (4.9) to an appropriate submanifold of \mathfrak{D} . Recall the decompositions (2.12), let us decompose the space \mathfrak{D} as

$$\mathfrak{D} = \mathfrak{D}_0 \oplus \mathfrak{D}_1, \tag{4.10}$$

where $\mathfrak{D}_{\nu} = \mathcal{D}_{\nu}^{-} \times \mathcal{D}_{\nu}^{+}$ for $\nu = 0, 1$. Since the subspaces \mathfrak{D}_{0} and \mathfrak{D}_{1} are dual to each other with respect to the inner product on \mathfrak{D} , then for any $\mathbf{A} \in \mathfrak{D}_{\nu}$ we have $T_{\mathbf{A}}^{*}\mathfrak{D}_{\nu} = (\mathfrak{D}_{\nu})^{*} = \mathfrak{D}_{1-\nu}$ for $\nu = 0, 1$. It is straightforward to verify the following lemma.

Lemma 4.3 The Poisson structure (4.9) on \mathfrak{D} can be properly restricted to each of its subspaces \mathfrak{D}_0 and \mathfrak{D}_1 .

5 Bihamiltonian representation of the two-component BKP hierarchy

In this section, we are to find a submanifold of \mathfrak{D} where the Poisson pencil for the two-component BKP hierarchy lies, then after a further reduction of the Poisson structure constructed in last section we will express the hierarchy (3.10), (3.11) to the form of Hamiltonian equations.

Recall the operators $P \in \mathcal{D}^-$, $\hat{P} \in \mathcal{D}^+$ given in (3.5), we let

$$\mathbf{A} = (P^2 D^{-1}, D\hat{P}^2). \tag{5.1}$$

It is easy to see that $\mathbf{A} \in \mathfrak{D}_1$ (see (4.10)), and $\mathbf{A} = (A, \hat{A})$ has the following expression:

$$A = P^{2}D^{-1} = D + \sum_{i \ge 0} (v_{-i}D^{-2i-1} + f_{-i}D^{-2i-2}), \tag{5.2}$$

$$\hat{A} = D\hat{P}^2 = \rho D^{-1}\rho + \sum_{i>1} (\hat{v}_i D^{2i-1} + \hat{f}_i D^{2i-2}), \quad \rho = \hat{u}_{-1}.$$
 (5.3)

Denote $\mathbf{v} = (v_0, v_{-1}, \dots, \hat{v}_0, \hat{v}_1, \dots)$ with $\hat{v}_0 = \rho^2$. Then the coordinate \mathbf{v} is related to \mathbf{u} given in (3.9) by a Miura-type transformation, while f_{-i} and \hat{f}_i are linear functions of derivatives of \mathbf{v} determined by the symmetry property $(A^*, \hat{A}^*) = -(A, \hat{A})$. Hence the flows of the hierarchy (3.10), (3.11) can be described in the coordinate \mathbf{v} .

Given any local functional $F \in \mathcal{F}_0$ (remind the notation \mathcal{F}_0 in last section), its variational gradient with respect to \mathbf{A} , say $\delta F/\delta \mathbf{A}$, is defined to be $\mathbf{X} = (X, \hat{X}) \in \mathfrak{D}$ with

$$X = \frac{1}{2} \sum_{i \ge 0} \left(\frac{\delta F}{\delta v_{-i}(x)} D^{2i} + D^{2i} \frac{\delta F}{\delta v_{-i}(x)} \right), \tag{5.4}$$

$$\hat{X} = \frac{1}{2} \sum_{i>0} \left(\frac{\delta F}{\delta \hat{v}_i(x)} D^{-2i} + D^{-2i} \frac{\delta F}{\delta \hat{v}_i(x)} \right). \tag{5.5}$$

In a coordinate-free way, $\delta F/\delta \mathbf{A} = \mathbf{X}$ can be defined by

$$\delta F = \langle \mathbf{X}, \delta \mathbf{A} \rangle, \quad \mathbf{X} \in \mathfrak{D}_0.$$
 (5.6)

Note that in the latter definition, the variational gradient is determined up to a kernel part $\mathbf{Z} = (Z, \hat{Z}) \in \mathfrak{D}_0$ such that

$$Z_{+} = 0, \quad \hat{Z}_{-} = 0, \quad \hat{Z}_{+}(\rho) = 0.$$
 (5.7)

Let us consider the coset $(D,0) + \mathcal{U}$ consisting of operators of the form (5.1), namely,

$$\mathcal{U} = (\mathcal{D}_1^-)_- \times ((\mathcal{D}_1^+)_+ \times \mathcal{M}), \quad \mathcal{M} = \{\rho D^{-1}\rho \mid \rho \in \mathcal{A}\}.$$
 (5.8)

Here \mathcal{M} is regarded as a 1-dimensional manifold with coordinate ρ , and this manifold has tangent spaces of the form

$$T_{\rho}\mathcal{M} = \{\rho D^{-1}f + fD^{-1}\rho \mid f \in \mathcal{A}\}.$$

So the tangent bundle, denoted by $T\mathcal{U}$, of the coset $(D,0) + \mathcal{U}$ has fibers

$$T_{\mathbf{A}}\mathcal{U} = (\mathcal{D}_{1}^{-})_{-} \times ((\mathcal{D}_{1}^{+})_{+} \times T_{\rho}\mathcal{M}), \quad \mathbf{A} \in (D,0) + \mathcal{U}, \tag{5.9}$$

while the cotangent bundle $T^*\mathcal{U}$ of $(D,0) + \mathcal{U}$ is composed of

$$T_{\mathbf{A}}^* \mathcal{U} = (\mathcal{D}_0^-)_+ \times ((\mathcal{D}_0^+)_- \times T_\rho^* \mathcal{M}), \quad T_\rho^* \mathcal{M} = \mathcal{A}. \tag{5.10}$$

From (5.4), (5.5) one sees that $\delta F/\delta \mathbf{A} \in T_{\mathbf{A}}^* \mathcal{U}$ for any $F \in \mathcal{F}_0$.

Now we are ready to do the desired reduction of the Poisson structure.

Lemma 5.1 The map

$$\mathcal{P}: T^*\mathcal{U} \to T\mathcal{U} \tag{5.11}$$

defined by the formula (4.9) is a Poisson tensor on the coset $(D,0)+\mathcal{U}$ that consists of operators of the form (5.1).

Proof. We only need to show that the map defined by (4.9) admits the restriction to the coset $(D,0) + \mathcal{U}$, i.e., the following map is well defined:

$$\mathcal{P}_{\mathbf{A}}: T_{\mathbf{A}}^* \mathcal{U} \to T_{\mathbf{A}} \mathcal{U}, \quad \mathbf{A} \in (D, 0) + \mathcal{U}.$$
 (5.12)

Assume $\mathbf{X} = (X, \hat{X}) \in T_{\mathbf{A}}^* \mathcal{U} \subset \mathfrak{D}_0$. It follows from Lemma 4.3 that $\mathcal{P}_{\mathbf{A}}(X) \in \mathfrak{D}_1$. More precisely, the first component of $\mathcal{P}_{\mathbf{A}}(X)$ belongs to $(\mathcal{D}_1^-)_-$. On the other hand, for any $\hat{Y} \in (\mathcal{D}^+)_+$ we have

$$(\hat{A}\hat{Y} + \hat{Y}^*\hat{A})_- = (\rho D^{-1}\rho\hat{Y} + \hat{Y}^*\rho D^{-1}\rho)_-$$

$$= -(\hat{Y}^*\rho D^{-1}\rho)_-^* + \hat{Y}^*(\rho)D^{-1}\rho$$

$$= \rho D^{-1}\hat{Y}^*(\rho) + \hat{Y}^*(\rho)D^{-1}\rho \in T_\rho\mathcal{M},$$

then by taking $\hat{Y} = (\hat{X}\hat{A})_+, (AX)_+$ it follows that the second component of $\mathcal{P}_{\mathbf{A}}(\mathbf{X})$ lies in $(\mathcal{D}_1^+)_+ \times T_\rho \mathcal{M}$. Thus $\mathcal{P}_{\mathbf{A}}(\mathbf{X}) \in T_{\mathbf{A}}\mathcal{U}$, i.e., the map (5.12) is well defined. The lemma is proved.

Remark 5.2 The proof of this lemma is the simplest case of the Dirac reduction procedure for Poisson tensors, see e.g. [20]. In fact, one can express the manifolds \mathfrak{D}_1 and \mathfrak{D}_1^* as

$$\mathfrak{D}_1 = \mathcal{U} \times \mathcal{V} = T_{\mathbf{A}} \mathcal{U} \times \mathcal{V}_{\mathbf{A}}, \quad \mathfrak{D}_1^* = \mathfrak{D}_0 = T_{\mathbf{A}}^* \mathcal{U} \times \mathcal{V}_{\mathbf{A}}^*, \tag{5.13}$$

where

$$\mathcal{V} = \mathcal{V}_{\mathbf{A}} = (\mathcal{D}_{1}^{-})_{+} \times \mathcal{N}, \quad \mathcal{N} = \{ X \in (\mathcal{D}_{1}^{+})_{-} \mid \text{res} X = 0 \},$$
$$\mathcal{V}_{\mathbf{A}}^{*} = (\mathcal{D}_{0}^{-})_{-} \times (T_{o}^{*})^{\perp} \mathcal{M}, \quad (T_{o}^{*})^{\perp} \mathcal{M} = \{ \hat{Y} \in (\mathcal{D}_{0}^{+})_{+} \mid \hat{Y}(\rho) = 0 \}.$$

Similar as the proof of Lemma 5.1, one can show that the map

$$\mathcal{P}_{\mathbf{A}} = \begin{pmatrix} \mathcal{P}_{\mathbf{A}}^{\mathcal{U}\mathcal{U}} & \mathcal{P}_{\mathbf{A}}^{\mathcal{U}\mathcal{V}} \\ \mathcal{P}_{\mathbf{A}}^{\mathcal{V}\mathcal{U}} & \mathcal{P}_{\mathbf{A}}^{\mathcal{V}\mathcal{V}} \end{pmatrix} : T_{\mathbf{A}}^{*}\mathcal{U} \times \mathcal{V}_{\mathbf{A}}^{*} \to T_{\mathbf{A}}\mathcal{U} \times \mathcal{V}_{\mathbf{A}}$$

defined by (4.9) is diagonal. Hence from Lemma 4.3 it follows that the map (4.9) gives a Poisson tensor on the coset $(D,0) + \mathcal{U} \subset \mathcal{D}_1$.

Lemma 5.3 On the coset $(D,0) + \mathcal{U}$ there are two compatible Poisson tensors defined by the following formulae:

$$\mathcal{P}_{1}(X,\hat{X}) = (A(XD^{-1})_{-} + D^{-1}(XA)_{-} - (D^{-1}X)_{-}A - (AX)_{-}D^{-1} - A(D\hat{X})_{-} - D^{-1}(\hat{A}\hat{X})_{-} + (\hat{X}D)_{-}A + (\hat{X}\hat{A})_{-}D^{-1},$$

$$\hat{A}(\hat{X}D)_{+} + D(\hat{X}\hat{A})_{+} - (D\hat{X})_{+}\hat{A} - (\hat{A}\hat{X})_{+}D - \hat{A}(D^{-1}X)_{+} - D(AX)_{+} + (XD^{-1})_{+}\hat{A} + (XA)_{+}D),$$

$$(5.14)$$

$$\mathcal{P}_2(X,\hat{X}) = (A(XA)_- - (AX)_- A - A(\hat{A}\hat{X})_- + (\hat{X}\hat{A})_- A,$$
$$\hat{A}(\hat{X}\hat{A})_+ - (\hat{A}\hat{X})_+ \hat{A} - \hat{A}(AX)_+ + (XA)_+ \hat{A})$$
(5.15)

with $(X, \hat{X}) \in T_{\mathbf{A}}^* \mathcal{U}$ at any point $\mathbf{A} = (A, \hat{A}) \in (D, 0) + \mathcal{U}$.

Proof. Lemma 5.1 shows that \mathcal{P}_2 is a Poisson tensor on the coset $(D,0) + \mathcal{U}$. Introduce a shift transformation on $(D,0) + \mathcal{U}$ as

$$\mathscr{S}: (A, \hat{A}) \mapsto (A + sD^{-1}, \hat{A} + sD)$$

with s being a parameter. Then the push-forward of the Poisson tensor \mathcal{P}_2 reads

$$(\mathscr{S}_* \mathcal{P}_2)(X, \hat{X}) = \mathcal{P}_2(X, \hat{X}) + s\mathcal{P}_1(X, \hat{X}) + s^2 \mathcal{P}_0(X, \hat{X}), \tag{5.16}$$

where

$$\mathcal{P}_0(X,\hat{X}) = (D^{-1}(XD^{-1})_- - (D^{-1}X)_- D^{-1} - D^{-1}(D\hat{X})_- + (\hat{X}D)_- D^{-1},$$

$$D(\hat{X}D)_+ - (D\hat{X})_+ D - D(D^{-1}X)_+ + (XD^{-1})_+ D).$$

By virtue of the symmetry property $(X^*, \hat{X}^*) = (X, \hat{X})$ that yields the formulae

$$(XD^{-1})_{\pm} = X_{\pm}D^{-1} \mp X_{+}(1)D^{-1},$$

$$(D^{-1}X)_{\pm} = D^{-1}X_{\pm} \mp D^{-1} \cdot X_{+}(1),$$

$$(D\hat{X})_{\pm} = D\hat{X}_{\pm}, \quad (\hat{X}D)_{\pm} = \hat{X}_{\pm}D,$$

one can check $\mathcal{P}_0(X, \hat{X}) = 0$. Hence the expansion (5.16) implies that \mathcal{P}_1 is a Poisson tensor that is compatible with \mathcal{P}_2 . The lemma is proved.

Let $\{\cdot,\cdot\}_{1,2}$ denote the Poisson brackets given in (4.8) with Poisson tensors being $\mathcal{P}_{1,2}$ respectively. We arrive at the main result of this article.

Theorem 5.4 The two-component BKP hierarchy (3.10), (3.11) can be expressed in the following bihamiltonian recursion form

$$\frac{\partial F}{\partial t_k} = \{F, H_{k+2}\}_1(\mathbf{A}) = \{F, H_k\}_2(\mathbf{A}),\tag{5.17}$$

$$\frac{\partial F}{\partial \hat{t}_k} = \{F, \hat{H}_{k+2}\}_1(\mathbf{A}) = \{F, \hat{H}_k\}_2(\mathbf{A})$$
 (5.18)

with $k \in \mathbb{Z}_+^{\text{odd}}$, where $F \in \mathcal{F}_0$, $\mathbf{A} = (P^2D^{-1}, D\hat{P}^2)$ as given in (5.1), and the Hamiltonians are

$$H_k = \frac{2}{k} \langle P^k \rangle, \quad \hat{H}_k = -\frac{2}{k} \langle \hat{P}^k \rangle, \quad k \in \mathbb{Z}_+^{\text{odd}}.$$
 (5.19)

Proof. First let us compute the variational gradients of the Hamiltonian functionals. Since

$$\delta H_k = \langle P^{k-2}, \delta P^2 \rangle = \langle DP^{k-2}, \delta(P^2D^{-1}) \rangle = \langle (DP^{k-2}, 0), \delta \mathbf{A} \rangle$$

and similarly

$$\delta \hat{H}_k = \langle (0, -\hat{P}^{k-2}D^{-1}), \delta \mathbf{A} \rangle$$

then up to kernel parts given in (5.7) we have the variational gradients of the Hamiltonians:

$$\frac{\delta H_k}{\delta \mathbf{A}} = (DP^{k-2}, 0), \quad \frac{\delta \hat{H}_k}{\delta \mathbf{A}} = (0, -\hat{P}^{k-2}D^{-1})$$
 (5.20)

One can easily see that different choices of the kernel parts do not change the definition of the Poisson tensors $\mathcal{P}_{1,2}$.

According to the flows (3.10), (3.11) one has

$$\frac{\partial \mathbf{A}}{\partial t_k} = \left([(P^k)_+, P^2] D^{-1}, D[(P^k)_+, \hat{P}^2] \right).$$

Note that

$$\frac{\partial F}{\partial t_k} = \left\langle \frac{\delta F}{\delta \mathbf{A}}, \frac{\partial \mathbf{A}}{\partial t_k} \right\rangle,$$

then to show (5.17) we only need to verify the equations

$$\frac{\partial \mathbf{A}}{\partial t_k} = \mathcal{P}_1 \left(\frac{\delta H_{k+2}}{\delta \mathbf{A}} \right) = \mathcal{P}_2 \left(\frac{\delta H_k}{\delta \mathbf{A}} \right). \tag{5.21}$$

The verification is straightforward by substituting (5.20) into (5.14), (5.15) with the help of the following formulae induced from (3.8):

$$(DP^kD^{-1})_{\pm} = D(P^k)_{\pm}D^{-1}, \quad (D\hat{P}^kD^{-1})_{\pm} = D(\hat{P}^k)_{\pm}D^{-1}, \quad k \in \mathbb{Z}_+^{\text{odd}}.$$

The equations (5.18) can be checked similarly. The theorem is proved. \Box

This theorem implies that the tau function (3.12) of the two-component BKP hierarchy is defined from the tau-symmetry of Hamiltonian densities [11] (up to the signs of \hat{H}_k).

Remark 5.5 One can also construct Hamiltonian structures of the two-component BKP hierarchy by reducing the linear and the cubic Poisson brackets induced from the R-matrix mentioned in last section. However, from these brackets we have not found bihamiltonian recursion relations like (5.17), (5.18).

6 Dispersionless limit of the bihamiltonian structure

Let us compute the leading term of the bihamiltonian structure in (5.17), (5.18) of the two-component BKP hierarchy, which would make sense in studying the corresponding Frobenius manifold if there be.

First we replace the pseudo-differential operators by Laurent series of symbols. In the dispersionless case, the operator $\mathbf{A} = (P^2D^{-1}, D\hat{P}^2)$ becomes

$$(a(z), \hat{a}(z)) = \left(z + \sum_{i \ge 0} v_{-i} z^{-2i-1}, \sum_{i \ge 0} \hat{v}_i z^{2i-1}\right), \tag{6.1}$$

and the coordinate-type local functionals $v_{-i}(y)$, $\hat{v}_j(y)$ have variational gradients $(z^{2i}\delta(x-y),0)$, $(0,z^{-2j}\delta(x-y))$ respectively. Substituting these Laurent series into the Poisson brackets defined by the formulae (4.8), (5.14), (5.15), we obtain the following result.

For the convenience of expression we set $v_1 = 1$, $v_i = 0$ when $i \ge 2$, and $\hat{v}_i = 0$ when $j \le -1$.

i) The first bracket: for $i, j \geq 0$,

$$\{v_{-i}(x), v_{-j}(y)\}_{1}^{[0]} = (1 - \delta_{i0} - \delta_{j0}) (2(i+j-1)v_{-i-j+1}(x) \delta'(x-y) + (2j-1)v'_{-i-j+1}(x) \delta(x-y)),$$
(6.2)

$$\{\hat{v}_i(x), \hat{v}_j(y)\}_1^{[0]} = -(1 - \delta_{i0} - \delta_{j0}) (2(i+j-1)\hat{v}_{i+j}(x) \, \delta'(x-y) + (2j-1)\hat{v}'_{i+j}(x) \, \delta(x-y)),$$

$$(6.3)$$

$$\begin{aligned}
\{v_{-i}(x), \hat{v}_{j}(y)\}_{1}^{[0]} \\
&= 2(i-j) \left((1-\delta_{j0})v_{j-i}(x) + (1-\delta_{i0})\hat{v}_{j-i+1}(x) \right) \delta'(x-y) \\
&- (2j-1) \left((1-\delta_{j0})v'_{j-i}(x) + (1-\delta_{i0})\hat{v}'_{j-i+1}(x) \right) \delta(x-y).
\end{aligned} (6.4)$$

ii) The second bracket: for $i, j \geq 0$,

$$\begin{aligned} \{v_{-i}(x), v_{-j}(y)\}_{2}^{[0]} \\ &= \sum_{r=-1}^{i-1} \left(2(i+j-2r-1)v_{-r}(x) \, v_{-i-j+r+1}(x) \, \delta'(x-y) \right. \\ &+ (2j-2r-1)v_{-r}(x) \, v'_{-i-j+r+1}(x) \, \delta(x-y) \\ &+ (2i-2r-1)v'_{-r}(x) \, v_{-i-j+r+1}(x) \, \delta(x-y) \right), \end{aligned} \tag{6.5}$$

$$\{\hat{v}_i(x), \hat{v}_j(y)\}_2^{[0]} = -\sum_{r=0}^i \left(2(i+j-2r+1)\hat{v}_r(x)\,\hat{v}_{i+j-r+1}(x)\,\delta'(x-y)\right)$$

$$+ (2j - 2r + 1)\hat{v}_{r}(x) \,\hat{v}'_{i+j-r+1}(x) \,\delta(x - y)$$

$$+ (2i - 2r + 1)\hat{v}'_{r}(x) \,\hat{v}_{i+j-r+1}(x) \,\delta(x - y) \Big), \qquad (6.6)$$

$$\{v_{-i}(x), \hat{v}_{j}(y)\}_{2}^{[0]}$$

$$= \sum_{r=\max\{-1, i-j-1\}}^{i-1} \Big(2(i-j)v_{-r}(x) \,\hat{v}_{-i+j+r+1}(x) \,\delta'(x - y)$$

$$+ (2r - 2j + 1)v_{-r}(x) \,\hat{v}'_{-i+j+r+1}(x) \,\delta(x - y)$$

$$+ (2r - 2i + 1)v'_{-r}(x) \,\hat{v}_{-i+j+r+1}(x) \,\delta(x - y) \Big). \qquad (6.7)$$

7 Concluding remarks

Based on the Lax pair representation (3.10), (3.11) of the two-component BKP hierarchy, we obtain a bihamiltonian structure of this hierarchy. Our method in the construction of the Poisson brackets is to employ the standard R-matrix formalism, which is analogous to that for the two-dimensional Toda hierarchy [1]. In comparison with the two-dimensional Toda hierarchy, we expect that there would be an infinite-dimensional Frobenius manifold underlying the two-component BKP hierarchy.

As shown in [19], the two-component BKP hierarchy (3.10), (3.11) is reduced to the Drinfeld-Sokolov hierarchy of type $(D_n^{(1)}, c_0)$ under the constraint $P^{2n-2} = \hat{P}^2$. Whether such a constraint induces a reduction of the bihamiltonian structure is unclear yet. We hope that considering this example would help to understand the relations between Frobenius manifolds of infinite and finite dimensions.

Acknowledgments. The authors thank Si-Qi Liu and Youjin Zhang for their advice, and thank Yang Shi for her comment. They are also grateful to the referee for the helpful suggestions. This work is partially supported by the National Basic Research Program of China (973 Program) No.2007CB814800 and the NSFC No.10801084.

References

- [1] Carlet, G. The Hamiltonian structures of the two-dimensional Toda lattice and *R*-matrices. Lett. Math. Phys. 71 (2005), no. 3, 209–226.
- [2] Carlet, G.; Dubrovin, B.; Mertens, L.P. Infinite-dimensional Frobenius manifolds for 2+1 integrable systems, preprint arXiv: math-ph/0902.1245v1.

- [3] Chen, Y.T.; Tu, M.H. On kernel formulas and dispersionless Hirota equations of the extended dispersionless BKP hierarchy. J. Math. Phys. 47 (2006), no. 10, 102702, 19 pp.
- [4] Date, E.; Jimbo, M.; Kashiwara, M.; Miwa, T. Transformation groups for soliton equations. IV. A new hierarchy of soliton equations of KP-type. Phys. D 4 (1981/82), no. 3, 343–365.
- [5] Date, E.; Jimbo, M.; Kashiwara, M.; Miwa, T. Transformation groups for soliton equations. Euclidean Lie algebras and reduction of the KP hierarchy. Publ. Res. Inst. Math. Sci. 18 (1982), no. 3, 1077–1110.
- [6] Date, E.; Kashiwara, M.; Jimbo, M.; Miwa, T. Transformation groups for soliton equations. Nonlinear integrable systems—classical theory and quantum theory (Kyoto, 1981), 39–119, World Sci. Publishing, Singapore, 1983.
- [7] Date, E.; Kashiwara, M.; Miwa, T. Transformation groups for soliton equations. II. Vertex operators and τ functions. Proc. Japan Acad. Ser. A Math. Sci. 57 (1981), no. 8, 387–392.
- [8] Dickey, L.A. Soliton equations and Hamiltonian systems. Second edition. Advanced Series in Mathematical Physics, 26. World Scientific Publishing Co., Inc., River Edge, NJ, 2003.
- [9] Drinfeld, V.G.; Sokolov, V.V. Lie algebras and equations of Kortewegde Vries type. (Russian) Current problems in mathematics, Vol. 24, 81– 180, Itogi Nauki i Tekhniki, Akad. Nauk SSSR, Vsesoyuz. Inst. Nauchn. i Tekhn. Inform., Moscow, 1984.
- [10] Dubrovin, B. Geometry of 2D topological field theories. Integrable systems and quantum groups (Montecatini Terme, 1993), 120–348, Lecture Notes in Math., 1620, Springer, Berlin, 1996.
- [11] Dubrovin, B.; Zhang, Y. Normal forms of integrable PDEs, Frobenius manifolds and Gromov-Witten invariants, preprint arXiv: math.DG/0108160.
- [12] Frenkel, E.; Givental, A.; Milanov, T. Soliton equations, vertex operators, and simple singularities, preprint arXiv: math.QA/0909.4032v1.
- [13] Givental, A.; Milanov, T. Simple singularities and integrable hierarchies. The breadth of symplectic and Poisson geometry, 173–201, Progr. Math., 232, Birkhäuser Boston, Boston, MA, 2005.
- [14] Jimbo, M.; Miwa, T. Solitons and infinite-dimensional Lie algebras. Publ. Res. Inst. Math. Sci. 19 (1983), no. 3, 943–1001.

- [15] Kac, V.G. Infinite-dimensional Lie algebras. Third edition. Cambridge University Press, Cambridge, 1990.
- [16] Kac, V.G.; Wakimoto, M. Exceptional hierarchies of soliton equations. Theta functions—Bowdoin 1987, Part 1 (Brunswick, ME, 1987), 191–237, Proc. Sympos. Pure Math., 49, Part 1, Amer. Math. Soc., Providence, RI, 1989.
- [17] ten Kroode, F.; van de Leur, J. Bosonic and fermionic realizations of the affine algebra \widehat{so}_{2n} . Comm. Algebra 20 (1992), no. 11, 3119–3162.
- [18] Li, L.C.; Parmentier, S. Nonlinear Poisson structures and r-matrices. Comm. Math. Phys. 125 (1989), no. 4, 545–563.
- [19] Liu, S.Q.; Wu C.Z.; Zhang Y. On the Drinfeld-Sokolov hierarchies of D type, preprint arXiv: nlin.SI/0912.5273v1.
- [20] Oevel, W.; Ragnisco, O. R-matrices and higher Poisson brackets for integrable systems. Phys. A 161 (1989), no. 1, 181–220.
- [21] Semenov-Tyan-Shanskii, M.A. What a classical r-matrix is. (Russian) Funktsional. Anal. i Prilozhen. 17 (1983), no. 4, 17–33. English translation: Functional Anal. Appl. 17 (1983), no. 4, 259–272.
- [22] Shiota, T. Prym varieties and soliton equations. Infinite-dimensional Lie algebras and groups (Luminy-Marseille, 1988), 407–448, Adv. Ser. Math. Phys., 7, World Sci. Publ., Teaneck, NJ, 1989.
- [23] Takasaki, K. Integrable hierarchy underlying topological Landau-Ginzburg models of *D*-type. Lett. Math. Phys. 29 (1993), no. 2, 111–121.
- [24] Takasaki, K. Dispersionless Hirota equations of two-component BKP hierarchy. SIGMA Symmetry Integrability Geom. Methods Appl. 2 (2006), Paper 057, 22 pp.
- [25] Ueno, K.; Takasaki, K. Toda lattice hierarchy. Group representations and systems of differential equations (Tokyo, 1982), 1–95, Adv. Stud. Pure Math., 4, North-Holland, Amsterdam, 1984.
- [26] Wu, C.Z. A remark on Kac-Wakimoto hierarchies of D-type. J. Phys. A: Math. Theor. 43 (2010), 035201.